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# Impact of Permanent Magnet Field on Inductance Variation of a PMLSM

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## Keywords

Linear drive, Modelling, Estimation technique, Synchronous motor, Harmonics.

## Abstract

Analytical models of Surface Mounted Permanent Magnet Linear Synchronous Motors are generally presented with constant inductances. However, the impact of powerful rare-earth PMs in the saturation phenomena cannot be neglected anymore. In this paper, a new non-linear model of PMLSM inductances is suggested. This model is defined as a function of the current value as well as of the magnet position. So, an analytical study, a finite-element analysis and experimental results are presented and confronted.

## 1. Introduction

Surface Mounted Permanent Magnet Linear Synchronous Machines (PMLSM) are widely used in the industrial framework [1]. These structures permit to achieve very high performances (power to weight ratio, high dynamics) [2]. The PMLSM's are often used as components for high-speed positioning systems, such as "Pick & Place" applications with a Gantry system [3]. Most of the time, the typical duty cycle uses high accelerations during short times, according to the energetic limitation induced by the  $I^2t$  value of the actuator. So the supply pulse current could be greater than the continuous rated current, which could imply saturation phenomena in the ferromagnetic sheets.

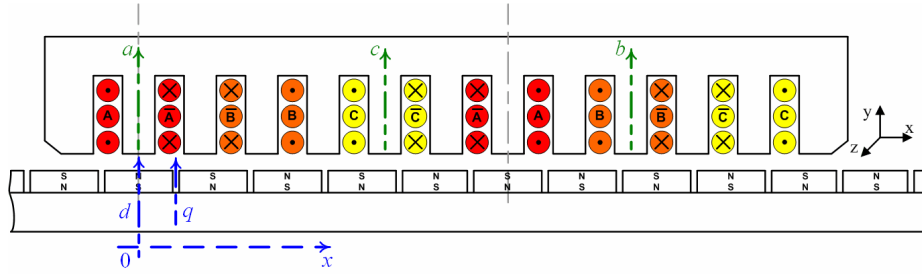
On the other hand, the PMLSM performances have been increased using rare-earth permanent magnets which have a very high residual magnetic flux density [4]. Besides, during the last decades, the price of PM such as Nd-Fe-B becomes more and more accessible [5]. The use of such PM leads to consider some modifications on the operating of the PMLSM. Indeed, due to the significant value of their residual flux density, the PM thickness needed is much smaller than in the case of the ferrite PM. This leads to a smaller equivalent air-gap with a high value of magnetic flux density, which is often close to the saturation level of the ferromagnetic sheets. Usually, for surface mounted PM synchronous motors, the inductances are assumed to be constant in classical models [6]. With the saturation phenomena, the inductances have to be taken as functions of both the current value and the magnet position.

First, an analytical model of the inductance variations is presented. The influence of both current and position variables is explained from an asymptotical point of view. The representation of the model is then obtained using the Causal Ordering Graph formalism. In a second part, the finite-element analysis using 2D-FEM software is presented. In the final part, experimental results are presented and compared with the finite element analysis. The measurement procedure is particularly detailed.

The investigated PMLSM is a flat ironcore LMD10-050 from Etel (Fig. 1). Table 1 lists some specifications of the PMLSM. This linear motor is a one pole pair structure with a pole pitch  $\tau_p = 16\text{mm}$ , and a force constant  $K_t$  of  $88.8\text{N/A}$ . The air gap thickness is about  $0.8\text{mm}$  and the Neodymium-Iron-Boron permanent magnets have a residual flux density of approximately  $1.23\text{Tesla}$ .

**Table 1: Specifications of the LMD10-050 [7]**

Rated Current	Maximal Current	Detent Force	Rated Thrust	Maximal Thrust
2Arms	7.9Arms	5N	130N	554N



**Fig. 1: Structure of the linear actuator LMD10-050 [8]**

## 2. Analytical model

The PMLSM studied is shown in figure 1. It is composed of a moving part (primary) and a fixed magnetic way (secondary). It is different from classical linear synchronous machines, as the magnet pole pitch  $\tau_p$  is smaller than the tooth pitch of the primary. Moreover, the ratio between both pole pitches has been designed so that the induced electromotive forces are sinusoidal functions of the secondary position [9], [10]. Besides, each armature phase of the machine is composed of two pole concentrated coils (Fig. 1). The electrical model of this PMLSM can be described by the classical relation [11]:

$$[v_{abc}] = R \cdot [i_{abc}] + \frac{d}{dt} [\phi_{abc}] \quad (1)$$

The three phase windings are star-connected, and the neutral-wire is accessible, but not connected. The winding distribution causes the mutual inductances to be negligible. Thus, the study of the non-linearity effects on a phase is independent from the other two. So, the flux of phase  $a$  is only a function of the current  $i_a$  and the position of the magnets  $x$ :

$$\frac{d\phi_a(x, i_a)}{dt} = \frac{\partial \phi_a}{\partial i_a} \cdot \frac{di_a}{dt} + \frac{\partial \phi_a}{\partial x} \cdot \frac{dx}{dt} \quad (2)$$

Then, we define the dynamic inductance  $L_{a\_dyn}$  as the partial derivative of the flux linkage with regard to the current, for a given position [12]:

$$L_{a\_dyn}(i_a, x) = \frac{\partial \phi_a(i_a, x)}{\partial i_a} \quad (3)$$

The general shape of the evolution of dynamic inductances can be deduced from a relatively simple reasoning based on the study of the variation of the magnetic permeability  $\mu$  of the material, according to the magnetic field  $H$ . In Fig. 2a, the simplified magnetic characteristic  $B(H)$  is composed of two different values of permeability:  $\mu_0$  in the saturated part, and  $\mu_0 \cdot \mu_r$  in the linear part of the material, where  $\mu_r = 4500$ . These values relate to the lower (saturated material) and the higher (linear zone) values, respectively  $L_{sat}$  and  $L_{lin}$ , on the dynamic inductance curve (Fig. 2b). The tooth magnetic field due to the magnets,  $H_m$ , varies with respect to the secondary position. It then constitutes a variable offset magnetic value. For a given position, the current  $I_{sat+}$ , which is necessary to reach the “positive saturation” level  $H_{sat}$  in the teeth, is proportional to  $\Delta H_+ = H_{sat} - H_m$ . For the same position, the current  $I_{sat-}$ , necessary to reach the “negative saturation” level  $-H_{sat}$ , is proportional to  $\Delta H_- = H_{sat} + H_m$ . These values,  $I_{sat+}$  and  $I_{sat-}$ , obviously depend on  $H_m$  and then on the secondary position, as shown in Fig. 2b. For the other two phases, the inductance variations can be deduced from those of phase *a* by a spatial phase shifting of  $\pm 2\tau_p/3$ .

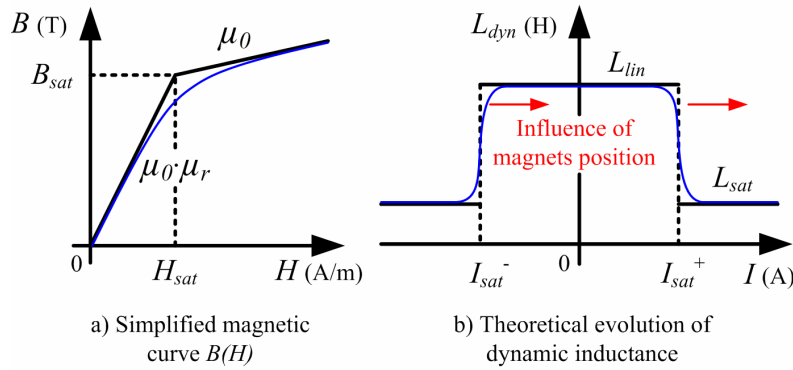


Fig. 2: Simplified  $B(H)$  and dynamic inductance curves

In order to analyse the influence of the dynamic inductance evolution on the PMLSM model, we use the Causal Ordering Graph (COG) formalism [13]. Indeed, the COG is based on a located energy representation with the theory of causal ordering [14]. It is undeniable that this approach is connected with that of links graphs such as the Bond Graphs [15]. But it differs from them by the analysis process, which is only based on integral causality. The COG is a tool that structures the synthesis of a state model, but that aims at maintaining, for a given system, a representation that stays as close as possible to its physical behaviour. In general, the expression of the transformation relations by means of the state equations is the best warranty against physical misinterpretation.

In the COG formalism, two complementary definitions of the causality are used. The corresponding COG symbols are presented in Fig.3:

- If an object accumulates information, the causality is internal. The output is necessarily a function of the energy state. The relation, thus oriented, is known as causal. The corresponding causal processor is given in Fig. 3a. Time and the initial state are implicit inputs and are not represented.
- On the contrary, if an object does not accumulate any information, causality is external. The output is thus an instantaneous function of the input. The relation, which is not oriented, is then known as rigid (Fig. 3b).

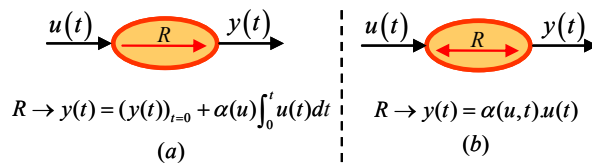


Fig.3: COG symbols:

(a) causal processor, (b) rigid processor.

The Causal Ordering Graph of the PMLSM is presented in Fig. 4. To facilitate the knowledge of the three phase modelling, this model is defined as a vector scheme. We notice in Eq.2 that the time derivative of the flux in the saturated case is considered as the sum of the flux partial derivative of the current function and flux partial derivative of the magnet position function. Using the COG, the calculation of these two partial derivatives is represented in the R6 processor. Indeed, in the saturated case, the flux has to be considered as a global function of the current and the magnet position. The  $\partial\Phi/\partial i$  is then used in the R2 processor, as the dynamic inductance to calculate the current value. The  $\partial\Phi/\partial x$  is used in the R5 processor, as the back electromotive force coefficient. The  $\partial\Phi/\partial x$  term is also used in the R4 processor for the electromechanical conversion. As the  $\partial\Phi/\partial x$  and the  $\partial\Phi/\partial i$  relations are time independent, the R6 processor can be considered as rigid. The R4 and the R5 processors are linked together as a gyrator [13]. That is classically known as a reversible electromechanical conversion.

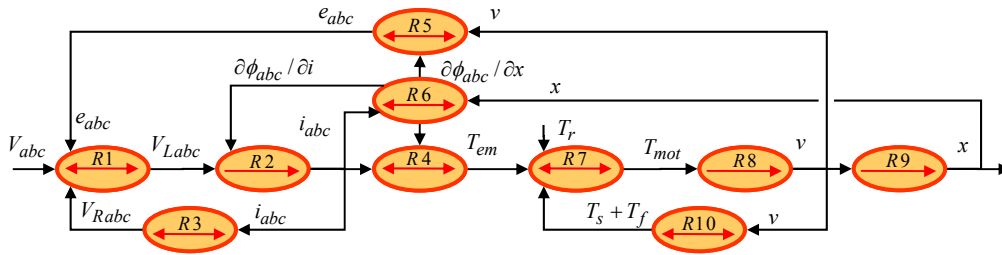


Fig. 4: COG representation of a PMLSM with saturation phenomena

Table II: Relation of the COG representation

Electrical Model:	Mechanical Model:
$R_1 \rightarrow V_{Labc} = V_{abc} - e_{abc} - V_{Rabc}$	$R_4 \rightarrow T_{em} = \sum_{k=a,b,c} \frac{\partial \phi_k(i_k, x)}{\partial x} \cdot i_k$
$R_2 \rightarrow \frac{\partial \phi_{abc}(i_{abc}, x)}{\partial i_{abc}} \cdot \frac{di_{abc}}{dt} = V_{Labc}$	$R_7 \rightarrow T_{mot} = T_{em} - T_r - T_s - T_f$
$R_3 \rightarrow V_{Rabc} = R \cdot i_{abc}, \quad R_5 \rightarrow e_{abc} = \frac{\partial \phi_{abc}(i_{abc}, x)}{\partial x} \cdot v$	$R_8 \rightarrow M \cdot \frac{dv}{dt} = T_{mot}, \quad R_9 \rightarrow \frac{dx}{dt} = v$
$R_6 \rightarrow \left. \frac{\partial \phi_{abc}(i_{abc}, x)}{\partial i_{abc}} = \frac{d\phi_{abc}(i_{abc}, x)}{di_{abc}} \right _{x=cst}$ $\left. \frac{\partial \phi_{abc}(i_{abc}, x)}{\partial x} = \frac{d\phi_{abc}(i_{abc}, x)}{dx} \right _{i_{abc}=cst}$	$R_{10} \rightarrow T_s + T_f = \begin{cases} Ts + f_v \cdot v, & si \ v > 0 \\ -Ts + f_v \cdot v, & si \ v < 0 \end{cases}$

### 3. Finite Element Analysis

To identify the dynamic inductances, we have solved a numerical model using the two-dimensional Finite Element (2D-FEM) calculation software called FEMM [16].

Because of the particular structure of the studied PMLSM, end effects of the finite length primary of the linear machine have to be taken into account. The finite element calculations are thus achieved with the complete actuator structure. On the other hand, the winding end effects cannot be taken into account with the 2D-FEM software. The non linear magnetic characteristic of the iron is taken into account with an average  $B(H)$  curve. Fig. 5 shows the primary of the LMD10-050 in two different positions with no supply current, according to the description of Fig. 1:

- The  $d$ -axis position ( $x = 0$ ), where one “North” magnet induces an important induction value in phase  $a$ . Position  $x$  is defined as the distance between the  $a$ -axis of the primary and the  $d$ -axis of the secondary.
- The  $q$ -axis position ( $x = \tau_p/2$ ), where phase  $a$  receives as much flux from one “South” magnet as from one “North” magnet.

Fig. 5 shows that the induction level in the teeth of phase  $a$ , for the  $d$ -axis position, is close to 1.6 Tesla. In consequence, the saturation phenomena in the ferromagnetic circuit can be reached more easily in the  $d$ -axis position than in the  $q$ -axis position.

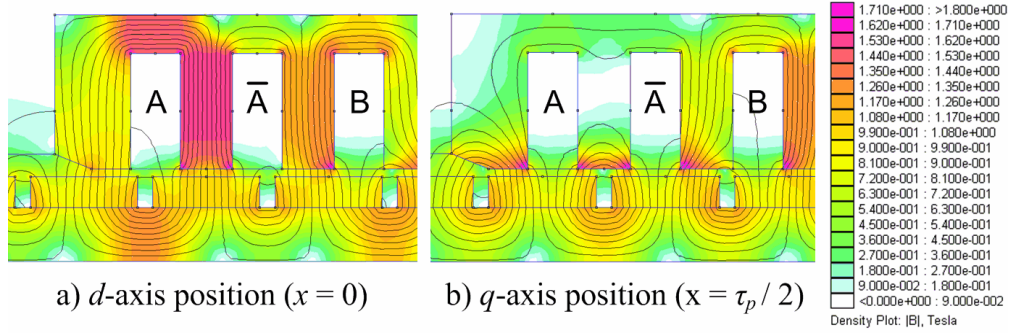


Fig. 5: Influence of permanent magnets position on the magnetic state of the structure with 2D-FEM calculations ( $i_a = i_b = i_c = 0$ )

For a given current value  $i_a$  and a given position  $x$ , the dynamic inductance  $L_{a\_dyn}(i_a, x)$  calculation is based on two flux values  $\Phi_{a1}(i_{a1}, x)$  and  $\Phi_{a2}(i_{a2}, x)$  obtained with a "little" current variation  $\Delta i_a = i_{a2} - i_{a1}$ . The dynamic inductance is thus calculated for the current value  $i_a = (i_{a1} + i_{a2}) / 2$ :

$$L_{a\_dyn}(i_a, x) = \frac{\phi_{a1} - \phi_{a2}}{i_{a1} - i_{a2}} = \frac{\Delta \phi_a}{\Delta i_a} \bigg|_{i_b = i_c = 0} \quad (4)$$

For accurate results, the current variation  $\Delta i_a$  has to be sufficiently weak to allow the linearization of the magnetic characteristic while limiting error on its slope. However, this variation must be sufficient to ensure the independence of results from the inaccuracy of this  $B(H)$  characteristic. Some experience has shown that a 0.01 Ampere  $\Delta i_a$  value is a good compromise.

Fig. 6 presents the finite element calculation results for the dynamic inductance of phase  $a$ .

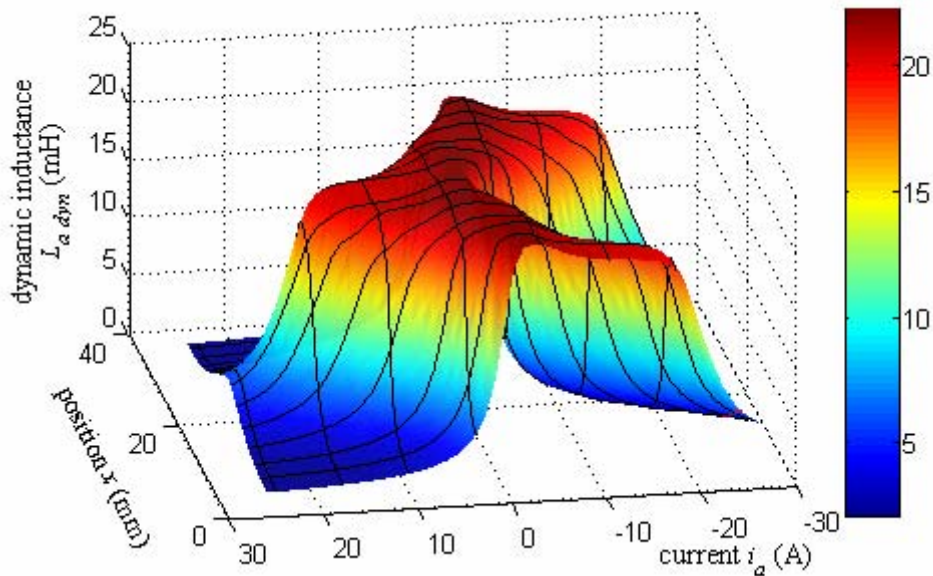


Fig. 6: Dynamic inductance of phase a using 2D-FEM



That confirms the illustration of Fig. 2 on the influence of the magnet position in the saturation phenomena. The calculations are performed with constant currents (magnetostatic case) which correspond here to an instantaneous current value. Additionally, the iso-position and iso-current curves are plotted over the 3D-Plot of the dynamical inductance in Fig. 6.

Fig. 7a shows the variation of the dynamic inductance  $L_{a\_dyn}$  as a function of the current  $i_a$  for 3 different positions of  $x$ . For  $x = \tau_p/2 = 8\text{mm}$ , the tooth of phase  $a$  is up to the middle of two opposite polarity magnets. So the saturation phenomenon is reduced. For  $x = 0\text{mm}$  the tooth of phase  $a$  is up to a North magnet. So the saturation phenomenon is maximal for positive current values. The impact of the permanent magnet could be considered in Fig. 7a, as a shifting of an equivalent sinusoidal current with a peak-to-peak value of about 18A. This peak-to-peak value is inversely proportional to the air-gap thickness.

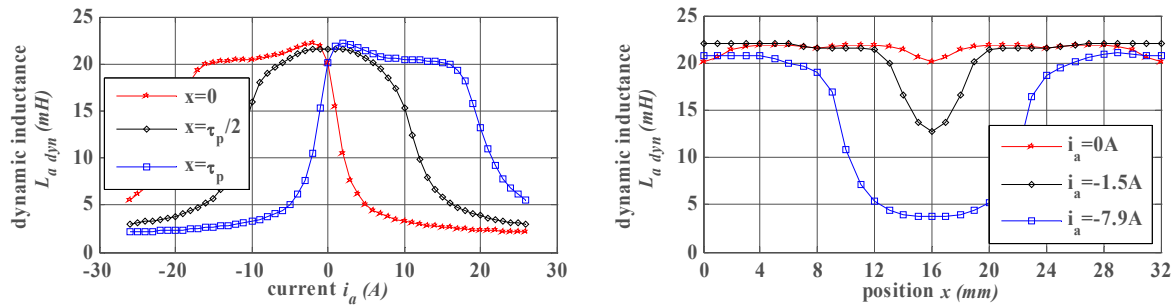


Fig. 7: Dynamic inductance  $L_{a\_dyn}$  as a function of:  
a) the current  $i_a$  for  $x = cst$       b) the position  $x$  and  $i_a = cst$

Fig. 7b shows the variation of the dynamic inductance  $L_{a\_dyn}$  as a function of position  $x$  for 3 different currents  $i_a$  (no current, rated current, peak current). For  $x = \tau_p = 16\text{mm}$ , phase  $a$  is up to a South magnet, so a negative current value increases the impact of the magnet position on the saturation phenomenon. With the maximal current value, the dynamic inductance is decreased by 80%.

## 4. Experimental validation

The proposed model is experimentally verified on a laboratory test system equipped with an ETEL LMD10-050 linear motor (Fig. 9). A Heidenhain exposed linear encoder with a grating period of  $20\mu\text{m}$  is used to detect the mover position. A dSPACE DS1005 real-time card is used for the identification procedures. The linear motor is star-connected, with an accessible neutral wire.

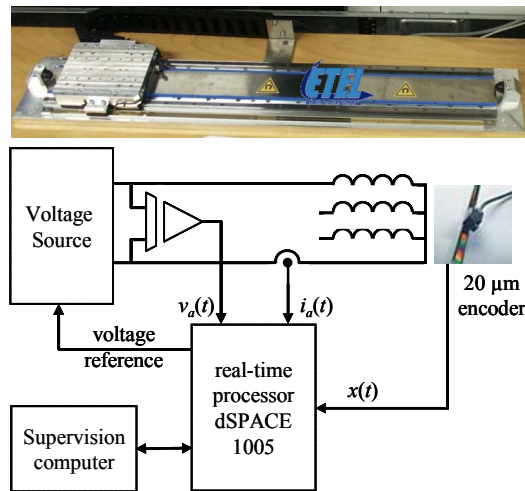


Fig. 9: Experimental test system with a LMD10-050 (ETEL) linear actuator

During the experiments, one phase of the PMLSM is supplied with a controlled sinusoidal voltage source (Fig. 9). According to Eq. 3, the dynamic inductance of phase  $a$  can be experimentally obtained while maintaining the primary in a given position. Fig. 10a shows the supplied voltage and the resulting non sinusoidal current, when  $x = \tau_p$ . During measurements, the current has been limited to the maximal admissible value of the peak current for the LMD10-050 (Table 1). Theoretically, it is possible to reach a higher current value using voltage pulses (close to the 26A of the finite elements results presented in Fig. 6).

The time variation of the flux in the studied phase is then obtained thanks to a time integration of the voltage and current measurements [17]:

$$\phi_a(t) = \int (v_a(t) - R \cdot i_a(t)) \cdot dt \quad (5)$$

The resulting non linear characteristic  $\Phi = f(i)$  is presented with position  $x = \tau_p$  (Fig. 10b). The effect of the non symmetrical saturation phenomena induced by the "South" magnet can be observed on both current and flux waveforms. The previously neglected hysteresis phenomenon induced by the ferromagnetic sheets can also be noticed. According to Eq. 3, the dynamic inductance of the considered phase is then deduced from an average non linear flux characteristic.

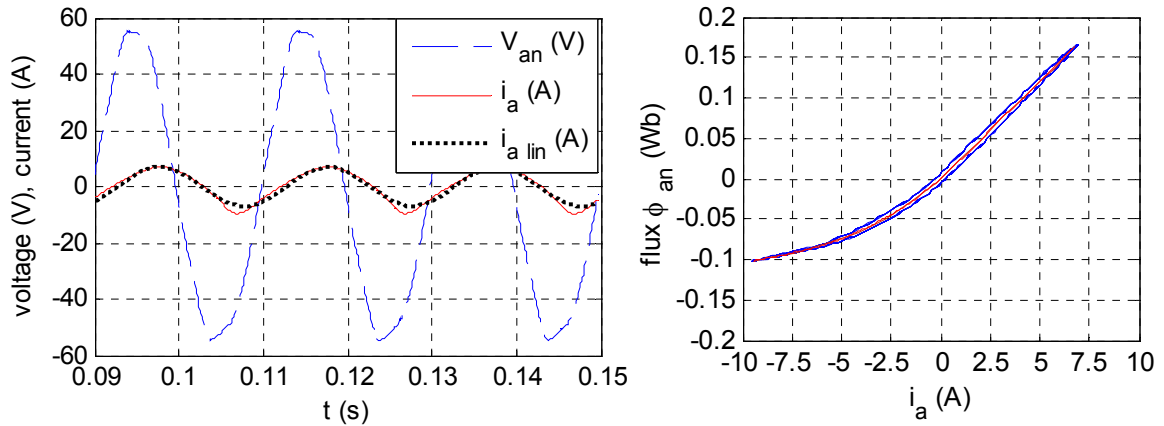


Fig. 10: Experimental identification:

a) voltage and current measurements

b) non linear flux characteristic  $\Phi=f(i)$

In Figure 11, the experimental dynamic inductance of phase  $a$  is compared to the FEM results when  $x = \tau_p$ . In this position, the tooth of phase  $a$  is up to a South magnet. The experimental identification results are about 10% higher than the finite element ones. This difference can be attributed to end-winding leakage inductances that are impossible to take into account with 2D FEM calculations. To achieve accurate 3D-FEM calculations, the number of elements has to be dramatically increased, leading to unacceptable computation time. Indeed, the required computation time to reach the 2D-FEM results is about 10 days on 4 PCs with 3GHz and 1GB RAMBUS.

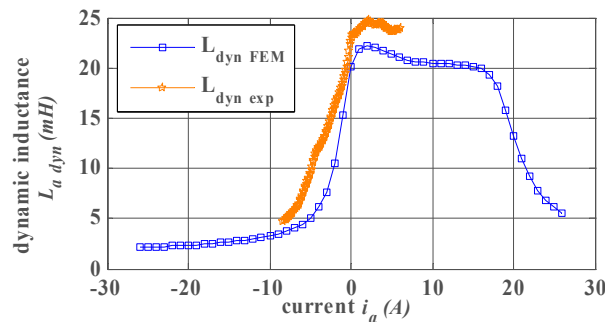


Fig. 11: Experimental dynamic inductance  $L_{a\_dyn}$  for  $x = \tau_p$



These end-winding effects are then approximately compensated in Fig. 12, increasing the FEM results (Fig. 6) by around 10% of the  $L_{a\_dyn}$  value. Experimental  $L_{a\_dyn}$  curves are given for three different positions ( $x = \tau_p / 3$ ,  $x = \tau_p$  and  $x = 5\tau_p / 3$ ) and are in good accordance with FEM results.

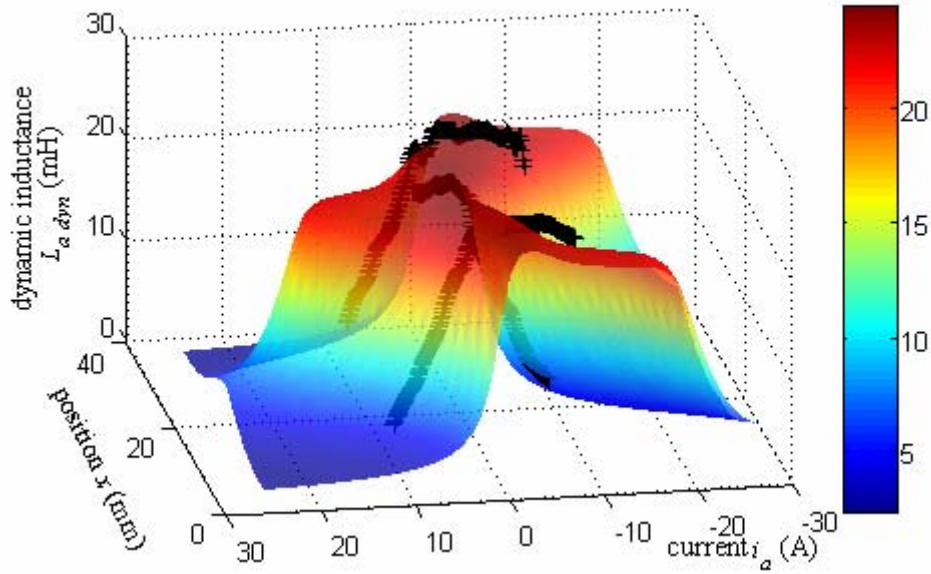


Fig. 12: Comparison of dynamic inductance  $L_{a\_dyn}$  for  $x = \tau_p$

## 5. Conclusion

In this paper, the inductance variations of a Surface Mounted Permanent Magnet Linear Synchronous Motor have been studied as a function of both the current and the magnet position. This non conventional approach is due to the saturation phenomena induced by the use of permanent magnets with very high residual magnetic flux density.

First, we have highlighted the origin of the saturation phenomena of the inductances. The proposed analytical model of the surface PMLSM is based on the dynamic inductances as a function of both the current and the position. Finite element analysis and experimental measurements validate the analytical model of the inductances proposed in this paper. Experimental results show a 10% error mainly caused by the neglected end-winding leakage inductances. As these effects cannot be integrated using 2D-FEM software, a future paper will describe a simplified 3D-FEM, regarding some computation time considerations.

Saturated inductances have many effects for the control design: non sinusoidal current, ripple thrust, peak thrust limitation. The study presented in this paper allows us to investigate further control design to reduce the saturation effects.

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